

# Virasoro type Lie algebras and deformations

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Received: 6 December 1995

**Abstract.** The first and second cohomologies of Cartan Type Lie algebras with coefficients in irreducible tensor modules are calculated. The space  $H^1(L, U)$  is interpreted as a space of deformations of  $(L, U)$ -modules.  $H^2(L, L) \neq 0$  if  $L = S_2, S_2^+$  or  $L = H_n, H_n^+$ . Lie algebra of divergenceless vector fields  $S_2^+$  has only one nontrivial local deformation. The two-sided simple hamiltonian algebra  $H_n$  has  $2n^2 + n$  new local deformations in addition to Moyal cocycle. The Lie algebras  $L = W_n (n > 3)$ ,  $S_{n-1} (n > 2)$ ,  $H_n (n > 1)$ ,  $K_{n+1} (n > 1)$  have 3, 1, 1, 3 nonisomorphic tensor modules with irreducible bases and nonzero 1-cohomologies; respectively, the corresponding numbers for 2-cohomologies are 9, 6, 7 and 9.

## 1 Introduction

The main applications of cohomologies of Lie algebras  $H^k(L, M)$  concern small  $k = 0, 1, 2, 3$ . The gauge group of cohomology introduced in physics by Faddeev [7] is related to the cohomology of the gauge Lie algebras [24]. Wess-Zumino functional is interpreted as a 1-cocycle of the space-time gauge group [23] and Schwinger term as a 2-cocycle of gauge group [25]. Abelian extensions appear in gauge theory and in gravitational theory. A 3-cocycles concern to the failure of the Jacobi identity in the presence of a Dirac monopole. Such fundamental results of the theory of finite-dimensional Lie algebras as Levi-Mal'cev theorem (any Lie algebra is a sum of a semisimple subalgebra and the radical) and Weyl theorem (any representation of a classical Lie algebra is a direct sum of irreducible representations) mean that, for a semisimple Lie algebra  $L$ , the first and second cohomologies with coefficients in irreducible modules are trivial.

In the nonclassical case both of these results are not true, but counterexamples to these facts also have important applications. Well known example of nonsplit extensions gives Virasoro algebra, i.e. a central extension of the

Lie algebra of vector fields on the circle [9]:

$$\begin{aligned} \text{Vir} = \{e_i, z : [e_i, e_j] &= (j - i)e_{i+j} + \delta(i + j = 0)(i^3 - i)z, \\ [e_i, z] &= 0, i, j \in \mathbb{Z}\} \end{aligned}$$

Indecomposable modules, outer derivations and generators of nilpotent algebras can be described in terms of first cohomologies. Second cohomologies are responsible for (non) split extensions and deformations [10]. Nonsplit extensions of Cartan Type Lie algebras called as Virasoro type Lie algebras. We interested in the following problem:

*For given Lie algebra  $L$  find irreducible modules  $M$  such that  $H^k(L, M) \neq 0$  and find basic cocycles*

As mentioned in [1] this problem for  $k = 2$  is originated from Cartan. We solve it for Cartan Type Lie algebras and their irreducible tensor modules in cases  $k = 1, 2$ .

Denote by  $\kappa_k(L)$  the number of irreducible modules  $M$  such that  $H^k(L, M) \neq 0$ . For any Lie algebra  $L$  of prime characteristic  $p$  the following inequality is true:  $0 < \kappa_k(L) < \infty$ ,  $0 < k < \dim L$  [3, 4]. For example,  $\kappa_2(sl_2) = 1$ ,  $p > 3$ . Some infinite-dimensional analoges of this result are also true for Lie algebras of Cartan type over complex numbers. From our results it follows that in the class of tensor modules,

$$\kappa_2(W_n) = 9 - \delta(n = 2, 3) - 4\delta(n = 1),$$

$$\kappa_2(S_{n-1}) = 6 - \delta(n = 3), n > 2,$$

$$\kappa_2(H_n) = 7 - 2\delta(n = 1),$$

$$\kappa_2(K_{n+1}) = 9,$$

$$\kappa_1(W_n) = \kappa_1(K_{n+1}) = 3,$$

$$\kappa_1(S_{n-1}) = \kappa_1(H_n) = 1$$

For example,  $W_1$  has five irreducible tensor modules with nonsplit extensions. One of them is trivial module and the corresponding extension is Virasoro algebra.

In our paper we consider two-sided Cartan type Lie algebras  $L$  (Lie algebras of vector fields in Laurent power series) and one-sided Cartan type Lie algebras  $L^+$  (Lie algebras of vector fields in formal power series). In Sect. 2

we give a constructive description of tensor modules for Cartan type Lie algebras. In Sect. 4 we describe all irreducible  $L_0^+$ -modules  $M_0$ , such that  $H^2(\mathcal{L}_0^+, M_0) \neq 0$ . As it turns out calculation of  $H^2(\mathcal{L}_0^+, M_0)$ , more precisely  $H^2(\mathcal{L}_1, \mathbb{C})$ , is equivalent to calculation of the following things:

- i) a space of defining relations of maximal graded nilpotent subalgebra  $\mathcal{L}_1^+$ ;
- ii) invariant bilinear forms  $L_1^+ \wedge L_1^+ \rightarrow M_0$  and sometimes  $L_k^+ \wedge L_q^+ \rightarrow M_0, k + q = 3, 4$ ;
- iii) all nonsplit extensions of  $L$  by irreducible  $(L, U)$ -modules.

Recall that  $L_0^+$  is isomorphic to one of the classical Lie algebras of types  $A_n, C_n$  or their split central extensions. We prove that all irreducible  $L_0^+$ -modules  $M_0$  with the property  $H^2(L, U \otimes M_0) \neq 0$  can be realised as subspaces of tensor modules of type  $(2, 4) (1, 3), (0, 2), (0, 1) (1, 2)$ , very like to the spaces of tensors of connection, curvature and etc. Here  $(s, k)$  denotes type of tensor,  $s$  is the number of contravariant components and  $k$  is the number of covariant components. For example,  $H^2(\mathcal{L}(S)_0^+, M_0) \neq 0$  exactly in the following five cases:

$$M_0 \cong \langle R_{ab} : R_{ab} = -R_{ba} \rangle \tag{0, 2},$$

$$M_0 \cong \langle R_{bcd}^a : R_{bcd}^a = R_{cbd}^a, R_{bcd}^a + R_{cdb}^a + R_{abc}^a = 0 \rangle \tag{1, 3},$$

$$M_0 \subseteq \langle R_{cdef}^{ab} : R_{cdef}^{ab} = -R_{efcd}^{ab} = R_{dcef}^{ab} = R_{cdfe}^{ab} = R_{cdef}^{ba} \rangle \tag{2, 4},$$

$$M_0 \subseteq \langle R_{cdef}^{ab} : R_{cdef}^{ab} = R_{dcef}^{ab} = R_{cedf}^{ab} = R_{cdfe}^{ab} = -R_{cdef}^{ba} \rangle \tag{2, 4},$$

$$M_0 \subseteq \langle R_{cdef}^{ab} : R_{cdef}^{ab} = R_{efcd}^{ab} = -R_{dcef}^{ab} = -R_{cdfe}^{ab} = -R_{cdef}^{ba} \rangle \tag{2, 4}$$

(here  $a, b, c, d, e, f = 0, 1, 2, \dots, n$ )

We would like to draw the attention to some corollaries of our result. The one-sided hamiltonian algebra  $H_n^+$  has only one nontrivial cocycle in the adjoint module, this cocycle is called the Moyal cocycle [17]. Corresponding local deformation can be prolonged to a global deformation and obtained deformation of Poisson bracket, called the Moyal bracket, is important in quantum mechanics. We prove that two-sided algebra  $[H_n, H_n]$  has, in addition to the Moyal cocycle,  $2n^2 + n$  more 2-cocycles in adjoint module. It is also interesting to note that Lie algebra of divergenceless vector fields on the 3-dimensional sphere  $S_2^+$  has exactly one nontrivial 2-cocycle in the adjoint module:

$$\psi(u\partial_i, v\partial_j) = \varepsilon_{abc}\partial_a\partial_j(u)\partial_b\partial_i(v)\partial_c,$$

where  $\varepsilon_{abc}$  is Levi-Chivita tensor. It seems to be that prolongation question of local deformations to global deformations is not so simple. In hamiltonian case  $H_n$ , for example, some of local deformations have obstructions.

The calculation of  $H^*(\mathcal{L}_0^+, M_0)$  may be useful for calculation of cohomologies or homologies of the nilpotent

subalgebra  $\mathcal{L}_1^+$  with coefficients in trivial module:

$$\bigoplus_{M_0} M_0 \otimes H^k(\mathcal{L}_0^+, L_0^+, M_0) \cong H^k(\mathcal{L}_1^+, \mathbb{C})$$

(here  $M_0$  runs through all the irreducible  $L_0^+$ -modules). Note that for Cartan Type Lie algebras  $L$  the cohomologies  $H^*(\mathcal{L}_1^+, \mathbb{C})$  was completely calculated only in case of  $L = W_1$ . [11]. At least eight proofs of this result are known, but it still remains one of the difficult results in cohomology theory of infinite-dimensional Lie algebras [8].

Another direction arising from our approach concerns defining relations of simple Lie algebras. As we mentioned above our problem is “almost” equivalent to the problem of calculating the second homologies of nilpotent subalgebra  $\mathcal{L}_1^+$ . The space  $H^2(\mathcal{L}_1^+, \mathbb{C})$  can be interpreted as a space of defining relations of  $\mathcal{L}_1^+$ . Defining relations of classical Lie algebras were found by Serre [22]. Nonclassical simple Lie algebras (Cartan Type Lie algebras) have big nilpotent subalgebras  $\mathcal{L}_1^+$  and  $H^2(\mathcal{L}_1^+, \mathbb{C})$  constitutes the main part of their defining relations. Calculations of  $H_2(\mathcal{L}_1^+, \mathbb{C})$  made jointly with Kerimbaev will be published elsewhere. Recall that 1-homologies of  $\mathcal{L}_1^+$  correspond to generators of  $\mathcal{L}_1^+$ . In characteristic zero generators were found by Gelfand and its collaborators [8], in characteristic  $p > 0$  it was calculated by Kostrikin and Shafarevich [13].

Many mathematicians and physicists were interested in generalizations of Virasoro algebras. Central extensions of two-sided Cartan type Lie algebras were described in [2]. In this paper proved that in the class of two-sided Cartan type Lie algebras only the following algebras have nonsplit central extensions:  $W_1$  (Virasoro algebra),  $S_n$  and  $H_n$ . From a physical viewpoint the problem of describing nonsplit extensions of Lie algebras of vector fields by modules of tensor fields was interested also in [15, 14, 16]. For  $L = W_n$  Larsson found two modules (differential 1-forms and 2-forms) having nontrivial 2-cocycles. For  $L = W_1, W_2$  and  $H_1$ , second cohomologies in all irreducible tensor modules was calculated in [4, 5] (actually in this paper the case of characteristic  $p > 0$  was considered, but the results contain the charactersitic 0 case as  $p \rightarrow \infty$ ). Nonsplit extensions of  $L = W_1$  independently described also in [21]. Some cocycles for  $L = W_n$  and  $L = H_n$  was also found in [20, 12, 18]

## 2 Preliminaries

All vector spaces are considered over  $\mathbb{C}$ . For a set of vectors  $\{u, v, \dots\}$  we will denote by  $\langle u, v, \dots \rangle$  its linear span. If  $\mathcal{A}$  is some statement, we will denote by  $\delta(\mathcal{A})$  its Kroneker symbol:  $\delta(\mathcal{A}) = 1$  if  $\mathcal{A}$  is true, and  $= 0$  if  $\mathcal{A}$  is false. Usually  $\delta(x = y)$  is denoted by  $\delta_{x,y}$ .

Let  $\Gamma_n$ , or more precisely  $\Gamma_I$  be a set of  $n$ -tuples  $\{\alpha = (\dots, \alpha_i, \dots), \alpha_i \in \mathbb{Z}, i \in I\}$ , where  $I$  is a set of indices and  $n = |I|$  is the number of its elements. Let  $\Gamma_n^+$  be the subset consisting of all  $\alpha$ , such that  $\alpha_i \geq 0$ . Let  $\mathbb{C}_I$  be the algebra of Laurent power series  $\mathbb{C}[x_i^\pm : i \in I] = \langle x^\alpha = \prod_{i \in I} x_i^{\alpha_i} : \alpha \in \Gamma_I \rangle$ . Instead of  $\mathbb{C}_I$  we shall write  $\mathbb{C}_{|I|}$  or simply  $U$  ( $I$  will be clear from the context). The Lie algebra

of General Type  $W_n$  is defined as the algebra of derivations of  $U$  with the usual commutator:

$$[u\partial_i, v\partial_j] = u\partial_i(v)\partial_j - v\partial_j(u)\partial_i, I = \{1, \dots, n\}.$$

Let

$$\omega_S = dx_0 \wedge dx_1 \wedge \dots \wedge dx_n, I = \{0, 1, \dots, n\}, \text{ (volume form)}$$

$$\omega_H = \sum_{i=1}^n dx_{-i} \wedge dx_i, I = \{\pm 1, \dots, \pm n\}, \text{ (hamiltonian form)}$$

$$\omega_K = dx_0 + \sum_{i=1}^n dx_{-i} \wedge dx_i, I = \{0, \pm 1, \dots, \pm n\}, \text{ (contact form)}$$

The Lie algebras of Special Type  $S_{n-1}$ , Hamiltonian Type  $H_n$  and Contact Type  $K_{n+1}$  are defined as the subalgebras of  $W_I$  saving the corresponding differential forms:

$$S_n = \{D \in W_{n+1}: D(\omega_H) = 0\},$$

$$H_n = \{D \in W_{2n}: D(\omega_H) = 0\},$$

$$K_{n+1} = \{D \in W_{2n+1}: D(\omega_K) = \omega_K\},$$

The Hamiltonian algebra can be defined as the vector space  $U$  with an even number variables with the Poisson bracket

$$\{u, v\} = \sum_i^n \text{sgn } i \partial_{-i}(u) \partial_i(v).$$

The Contact algebra is defined on vector space  $U$  with odd number variables as follows:

$$[u, v] = \partial_0(u)\Delta(v) - \partial_0(v)\Delta(u) + \sum_i^n \text{sgn } i \partial_{-i}(u) \partial_i(v),$$

where,  $\Delta(u) = (2 - \sum_{i \neq 0} x_i \partial_i)(u)$ . In general these algebras are not simple. To obtain simple ones we take second commutators and factorize by center (in hamiltonian case). Lower indices denote dimensions of standard tori. We endow  $U$  and its derivation algebras with gradings:

$$U = \bigoplus_k U_k, L = \bigoplus_k L_k, U_k U_q \subseteq U_{k+q}, [L_k, L_q] \subseteq L_{k+q}$$

$$U_k = \{x^\alpha: |\alpha| = \sum_{i \in I} \alpha_i = k\},$$

$$L = W_n, L_k = \{x^\alpha \partial_i: |\alpha| = k + 1, \},$$

$$L = H_n, L_k = \{x^\alpha: |\alpha| = k + 2\},$$

$$L = K_{n+1}, L_k = \left\{ x^\alpha: 2\alpha_0 + \sum_{i=1}^n \alpha_i = k + 2 \right\}$$

The filtrations corresponding to these gradins are:

$$L = \bigcup_k \mathcal{L}_k, \mathcal{L}_k = \bigoplus_{j \geq k} L_j, \mathcal{L}_k \supseteq \mathcal{L}_{k+1}, k \in \mathbb{Z}$$

Here gradations and filtrations run through all positive as well as neagive integers. This is why we call these algebras

two-sided Cartan type Lie algebras. They have subalgebras that we call one-sided Cartan Type Lie algebras, because their gradations and filtrations run integers from  $-2$  or  $-1$  only to positive parts. Namely, for the algebra  $U^+ = \mathbb{C}[[x_i: i \in I]]$  we introduce one-sided Cartan type Lie algebras as follows:

$$W_n^+ = \text{Der } U^+, S_n^+ = W_n^+ \cap S_n,$$

$$H_n^+, K_{n+1}^+ \text{ are subalgebras of } H_n, K_{n+1} \text{ spanned on } U^+$$

Then for  $L = W_n, S_n, H_n, K_{n+1}$  algebra  $L^+$  has gradation:

$$L^+ = \bigoplus_{k \geq -2} L_k^+, L_k^+ = L^+ \cap L_k$$

and filtration

$$\mathcal{L}_k^+ = \bigoplus_{j \geq k} L_j^+, L = \mathcal{L}_{-2}^+ \subseteq \mathcal{L}_{-1}^+ \subseteq \mathcal{L}_0^+ \subseteq \mathcal{L}_1^+ \dots$$

The subalgebra  $L_0^+$  is isomorphic to  $gl_n, sl_n, sp_n, sp_n \oplus \mathbb{C}$  for  $L = W_n, S_n, H_n, K_{n+1}$  respectively.

Let  $J(L)$  be a subalgebra  $L_{-1}^+$ , if  $L = W_n, S_{n-1}$  and  $L_{-1}^+ + L_{-2}^+$ , if  $L = H_n, K_{n+1}$ .

### 3 Tensor modules

Let  $Q$  be a Lie algebra of derivations of an associative commutative algebra  $V$ . We say that  $M$  is a  $(Q, V)$ -module if  $M$  has structures of module over the Lie algebra  $Q$  and over the associative algebra  $V$ , such that

$$D(vm) = D(v)m + vD(m)$$

for any  $D \in Q, v \in V$  and  $m \in M$ . For  $L = W_n, S_n, H_n, K_n, Q = L, L^+, V = U, U^+$  denote by  $\mathfrak{A}(V)$  the category of  $(Q, V)$ -modules. Let  $\mathfrak{A}_0$  be the category of  $L_0^+$ -modules. Make any  $M_0 \in \mathfrak{A}_0$  a module over  $\mathcal{L}_0^+$  by  $\mathcal{L}_1^+ M_0 = 0$ . Define a Functor

$$\mathfrak{A}_0 \rightarrow \mathfrak{A}(V), M_0 \mapsto V \otimes M_0$$

by the rule

$$D(v \otimes m) = D(v)m + \sum_a v E_a(D) \otimes a(m),$$

$$u(v \otimes m) = uv \otimes m, \forall u, v \in U, \forall D \in L, \forall m \in M_0,$$

where  $E_a: Q \rightarrow V$  are the linear maps invariant under  $L_1^+ + L_2^+$  constructed in [6] for any basic element  $a \in \mathcal{L}_0^+$ . We call the module  $V \otimes M_0$  the *tensor  $Q$ -module with base  $M_0$* . So, actions of  $L$  and its subalgebra  $L^+$  on tensor modules are given by the formulas

$$L = W_n, (u\partial_i)(v \otimes m) = u\partial_i(v) \otimes m + \sum_s \partial_s(u) \otimes x_s \partial_i(m),$$

$$L = S_n, (\partial_i(u)\partial_j - \partial_j(u)\partial_i)(v \otimes m)$$

$$= (\partial_i(u)\partial_j(v) - \partial_j(u)\partial_i(v)) \otimes m$$

$$+ \sum_s v \partial_s \partial_i(u) \otimes x_s \partial_j(m) - v \partial_s \partial_j(u) \otimes x_s \partial_i(m),$$

$$L = H_m, u(v \otimes m) = \{u, v\} \otimes m + (1/2) \sum_{i,j} u \partial_i \partial_j (v) \otimes x_i x_j (m),$$

$$L = K_{n+1}, u(v \otimes m)$$

$$= ([u, v] - 2\partial_0(u)v) \otimes m + v\partial_0(u) \otimes x_0(m) + (1/2) \sum_{i,j \neq 0} v(\partial_i + \text{sgn } i x_{-i} \partial_0)(\partial_j + \text{sgn } j x_{-j} \partial_0)(u) \otimes x_i x_j (m)$$

Notice that  $(L^+, U^+)$ -modules can be described in the language of induced or coinduced modules. For example,  $U^+ \otimes M_0 \cong \text{Ind}(\mathcal{L}_0, M_0)$ . But an  $(L, U)$ -module in general is not isomorphic either to an induced or a coinduced module.

For Cartan Type Lie algebra of toroidal dimension  $n$ ,  $(L, U)$ -modul  $M = U \otimes M_0$  has weight decomposition  $M = \bigoplus_{\alpha \in \Gamma_n} M_\alpha$  such that for any weight space  $M_\alpha = \{x^\alpha \otimes M_0\}$  we have  $\dim M_\alpha = \dim M_0$ . This property was used in [19] for constructing a  $W_n$ -module with weight subspaces of dimension  $n$  (take  $M_0$  as  $\Lambda^1$ ),  $n^2$  (take  $M_0$  as the adjoint module of  $gl_n$ ),  $n^3$ , etc.

#### 4 Cohomologies and interpretations

For a Lie algebra  $Q$  and a  $Q$ -module  $M$  denote by  $C^k(Q, M)$  the space of skewsymmetric polylinear maps with  $k$  arguments in  $Q$  and coefficients in  $M$ , if  $k > 0$ ,  $C^0(Q, M) = M$ , and  $C^k(Q, M) = 0$ , if  $k < 0$ . Let  $\varrho: Q \rightarrow \text{End} C^k(Q, M)$  be the standard representation of  $Q$ :

$$\varrho(D)\psi(D_1, \dots, D_k) = D(\psi(D_1, \dots, D_k)) + \sum_{i=1}^k (-1)^i \psi([D, D_i], D_1, \dots, \hat{D}_i, \dots, D_k)$$

and let  $\iota: C^k(Q, M) \rightarrow C^{k-1}(Q, M)$  be the inner product:

$$\iota(D)\psi(D_1, \dots, D_{k-1}) = \psi(D, D_1, \dots, D_{k-1})$$

(here  $\hat{D}$  means that the element  $D$  is omitted). For a subalgebra  $R$  of  $Q$  denote by  $C^k(Q, R, M)$  the subspace of  $C^k(Q, M)$  consisting of the maps  $\psi$  such that  $\iota(D)\psi = 0$  and  $\varrho(D)\psi = 0$  for any  $D \in R$ . Let  $C^*(Q, M) = \bigoplus_k C^k(Q, M)$ . In the cochain complex  $C^*(Q, M)$  the coboundary operator

$$d: C^k(Q, M) \rightarrow C^{k+1}(Q, M)$$

is defined by

$$d\psi(D_1, \dots, D_{k+1}) = \sum_{i < j} \psi([D_i, D_j], D_1, \dots, \hat{D}_i, \dots, \hat{D}_j, \dots, D_{k+1}) + \sum_{i=1}^{k+1} (-1)^i D_i \psi(D_1, \dots, \hat{D}_i, \dots, D_{k+1})$$

Let

$$Z^*(Q, M) = \text{Ker } d \text{ (space of cycles),}$$

$$B^*(Q, M) = \text{Im } d \text{ (space of coboundaries),}$$

$$H^*(Q, M) = Z^*(Q, M)/B^*(Q, M) \text{ (space of cohomologies).}$$

For example,

$$H^0(Q, M) = M^L = \{m \in M : D(m) = 0, \forall D \in Q\}$$

(submodule of invariants),

$$H^1(Q, \mathbb{C}) = Q/[Q, Q] \text{ (subspace of generators),}$$

$$H^1(Q, Q) \text{ (space of derivations),}$$

$H^2(L, L)$  is the space of local deformations (see [10]).

For Cartan Type Lie algebra  $L$  and its nilpotent subalgebra  $\mathcal{L}_1^+$  the space  $H^2(\mathcal{L}_1^+, \mathbb{C})$  can be interpreted as a space of defining relations of  $\mathcal{L}_1^+$  (see [8]). In next section we give an interpretation of  $H^1(L, U)$ ,  $H^1(L^+, U^+)$  as a spaces of deformations of representations.

For subalgebra  $R$  of  $Q$  relative cohomologies are defined as the cohomologies of subcomplex  $C^*(Q, R, M) = \bigoplus_k C^k(Q, R, M)$ .

Let  $L$  be a graded Lie algebra of Cartan Type:  $L = \bigoplus_k L_k$  and  $\mathcal{L}_0^+ = \bigoplus_{k>0} L_k^+$ ,  $M_0$ - $\mathcal{L}_0^+$ -module such that  $\mathcal{L}_1^+ M_0 = 0$ . The cochain complex  $C^*(\mathcal{L}_0^+, M_0)$  also has a gradation:

$$C^*(\mathcal{L}_0^+, M_0) = \bigoplus_{r \geq 0} C_r^*(\mathcal{L}_0^+, M_0),$$

$$C_r^*(\mathcal{L}_0^+, M_0) = \bigoplus_k C_r^k(\mathcal{L}_0^+, M_0),$$

$$C_r^k(\mathcal{L}_0^+, M_0) = \left\{ \psi \in C^k(\mathcal{L}_0^+, M_0) : \psi(D_1, \dots, D_k) = 0, \right. \\ \left. D_i \in L_{k_i}, \sum_i k_i \neq k \right\}.$$

Denote by  $H_r^*(\mathcal{L}_0^+, M_0)$  the cohomologies of the cochain subcomplex  $C_r^*(\mathcal{L}_0^+, M_0)$ .

For a Cartan Type Lie algebra  $L$  and  $Q = L, L^+, V = U, U^+$  the following map induces isomorphism of cochain complexes:

$$E: C^k(\mathcal{L}_0^+, M_0) \rightarrow C^k(Q, J(L), V \otimes M_0),$$

$$E\psi(D_1, \dots, D_k) = \sum_{a_1, \dots, a_k} E_{a_1}(D_1) \cdots E_{a_k}(D_k) \otimes \psi(a_1, \dots, a_k)$$

here  $a_1, \dots, a_k$  runs through the basic vectors of  $\mathcal{L}_0^+$ .

#### 5 Deformations of $(Q, V)$ -modules

Let  $M$  be an  $(Q, V)$ -module with action  $Q \times M \rightarrow M, (D, m) \mapsto D(m)$ . Let  $M$  be exact as a  $V$ -module:  $u(m) = 0, \forall m \in M \Rightarrow u = 0$ . Define on  $M_\lambda = M \otimes \mathbb{C}\{\lambda\}$  following Gerstenhaber a new structure of  $(Q, V)$ -module with action of  $Q$  given by a power series

$$D_\lambda(m) = D(m) + \lambda f_1(D)m + \lambda^2 f_2(D)m + \dots$$

We do not change the action of  $V$ . We shall say that  $M_\lambda$  is a deformation of the  $(Q, V)$ -module  $M$  and denote it by  $f = (f_1, f_2, \dots)$ , if these actions really give  $(Q, V)$ -module

structures. Deformations  $\mathbf{f}$  and  $\mathbf{g}$  are *equivalent* if the following diagramm is commutative

$$\begin{array}{ccc} M_\lambda & \xrightarrow{f} & M_\lambda \\ \downarrow \Phi & & \downarrow \Phi \\ M_\lambda & \xrightarrow{g} & M_\lambda \end{array}$$

under with

$$\Phi: M_\lambda \rightarrow M_\lambda, m \mapsto m + \lambda \phi_1(m) + \lambda^2 \phi_2(m) + \dots,$$

where  $\phi_1, \phi_2, \dots$  are linear operators on  $M$ . The condition that  $D \rightarrow D_\lambda$  is a representation of a Lie algebra is equivalent to

$$f_k \in Z^1(Q, V), \forall k > 0.$$

In particular,

$$f_1 \in Z^1(Q, V)$$

and for equivalent deformations  $\mathbf{f}, \mathbf{g}$ ,

$$f_1 - g_1 = d\phi \in B^1(Q, V).$$

Moreover any cocycle  $f \in Z^1(Q, V)$  as a local deformation  $f_1 = f$  can be prolonged to global deformation  $\mathbf{f}$  in a trivial way:  $f = (f_1, 0, \dots)$  is a deformation of  $M$ . So, we give an interpretation of the first space of cohomology  $H^1(Q, V)$  as a space of deformations of  $(Q, V)$ -modules.

It is easy to check that for a Cartan Type Lie algebra  $Q$  (one-sided or two-sided does not matter) and a two-sided module  $U$  the 1-cochains  $Sq_i \in C^1(Q, U)$ ,  $i \in I(L)$  defined by

$$Sq_i(D) = D(\ln x_i) : D \mapsto x_i^{-1} D(x_i),$$

are cocycles. The same is true for the following cochains in  $C^1(Q, V)$ :

$$Div: u \partial_i \mapsto \partial_i(u),$$

for  $Q = W_n, W_n^+, V = U, U^+$ ,

$$\Delta: u \mapsto \left(2 - \sum_i x_i \partial_i\right)(u), D_\theta: x^z \mapsto \delta_z = -\theta,$$

for  $L = H_n, H_n^+, V = U, U^+$ , (here  $\theta = (1, \dots, 1) \in \Gamma_{2n}$ )

$$\partial_0: u \mapsto \partial_0(u),$$

for  $L = K_{n+1}$ .

**Theorem 5.1.** Let  $Q = L, L^+$ , where  $L = W_n, S_{n-1}, H_n, K_{n+1}$ , and  $V = U, U^+$ , where  $U = \mathbb{C}[[x_i^\pm : i \in I = I(L)]]$ . Then

$$H^1(L, U) \cong H^1(L^+, U) \cong H^1(J(L), U) \oplus H^1(\mathcal{L}_0, \mathbb{C}),$$

$$H^1(L^+, U^+) \cong H^1(J(L), U^+) \oplus L_0^+ / [L_0^+, L_0^+],$$

$$H^1(J(L), V) \cong H^1(\Omega(V)) \cong \mathbb{C}^{|I(L)|},$$

if  $L \neq H_n$ , and

$$H^1(J(L), V) \cong H^1(\Omega(V)) \oplus \langle \Delta \rangle \cong \mathbb{C}^{2n+1},$$

if  $L = H_n$  where

$$H^1(\Omega(U)) \cong \langle d(\ln x_i) : i \in I(L) \rangle$$

$$\cong \langle Sq_i : i \in I(L) \rangle, H^1(\Omega(U^+)) = 0,$$

$$H^1(\mathcal{L}_0^+, \mathbb{C}) \cong \langle Div \rangle \cong \mathbb{C}, \text{ if } L = W_n, \cong \langle D_\theta \rangle \cong \mathbb{C}, \text{ if } L = H_n, \text{ and } \cong \langle \partial_0 \rangle \cong \mathbb{C}, \text{ if } L = K_{n+1}.$$

Let  $\Lambda^k = \langle dx_{i_1} \wedge \dots \wedge dx_{i_k} : i_1, \dots, i_k \in I \rangle$  be the  $L_0^+$ -module of  $k$ -differential forms,  $\Lambda^* = \bigoplus_k \Lambda^k$ , and  $\Omega^k(V) = V \otimes \Lambda^k$ ,  $\Omega^* = \bigoplus \Omega^k(V)$ . Endow  $\Omega^*(V)$  with a coboundary operator:

$$\Omega^k(V) \rightarrow \Omega^{k+1}(V),$$

$$d(v \otimes dx_{i_1} \wedge \dots \wedge dx_{i_k}) = \sum_{i \in I} \partial_i(v) \otimes dx_i \wedge dx_{i_1} \wedge \dots \wedge dx_{i_k}.$$

In this way we obtain the de Rham complex  $\Omega^*(V)$  with coefficients in  $V$ . Let, as usual,

$$Z^*(\Omega(V)) = \text{Kerd} \text{ (space of closed forms)}$$

$$B^*(\Omega(V)) = \text{Im} d \text{ (space of exact forms)}$$

and,

$$H^*(\Omega(V)) = Z^*(\Omega(V)) / B^*(\Omega(V)) \text{ (de Rham cohomologies).}$$

**Proposition 5.2.**  $H^*(U) \cong \Lambda^* = \langle d(\ln x_{i_1}) \wedge \dots \wedge d(\ln x_{i_k}) : i_1, \dots, i_k \in I \rangle$ ,  $H^*(U^+) \cong 0$ .

In particular,

$$\langle Sq_i : i \in I \rangle \cong H^1(\Omega(U)), Sq_i \mapsto d(\ln x_i).$$

Recall that any irreducible  $L_0^+$ -module  $M_0$  is uniquely determined by its highest weight  $\pi$ , and conformal weight  $\lambda$  (if  $L_0^+$  is simple, then  $\lambda = 0$ ).  $L_0^+$  has 1-dimensional center in the following cases:

$$L = W_n, L_0^+ = \langle x_i \partial_j : i, j = 1, \dots, n \rangle \cong gl_n,$$

$$(x_i \partial_j)_\lambda(m) = x_i \partial_j(m) + (1 - \lambda) \delta(i = j)m, m \in M_0$$

and

$$L = K_{n+1},$$

$$L_0^+ = \langle x_i x_j : i, j = \pm 1, \dots, \pm n \rangle \oplus \langle x_0 \rangle$$

$$\cong sp_n \oplus \mathbb{C},$$

$$u_\lambda(m) = u(m) + (\lambda - 2) \partial_0(u)m, u \in L_0^+, m \in M_0.$$

We will write  $M = R(\pi)$  or  $M = R(\pi, \lambda)$ . Denote by  $\pi_i$ ,  $1 \leq i \leq \dim T$ , the fundamental weights of  $L_0^+$ . Any highest weight can be represented as a linear combination of fundamental weights with nonnegative integer coefficients  $\pi = \sum_i l_i \pi_i$ . It is well known that any irreducible module  $M_0 = R(\pi)$  is isomorphic to a submodule (or factor-module) of some module  $\hat{M}_0$ , obtained from tensor ( $\otimes$ ), exterior ( $\wedge$ ) and symmetric ( $\circ$ ) products of fundamental representations  $R(\pi_i)$ :

$$\hat{M}_0 = \otimes_i S^{l_i}(\Lambda^i).$$

For tensors  $a$  and  $b$  the following notations are used:  $a \wedge b = a \otimes b - b \otimes a$ ,  $a \circ b = a \otimes b + b \otimes a$ . Let  $\Lambda^k = \langle dx_{i_1} \wedge \dots \wedge dx_{i_k} \rangle$  be the  $k$ -th exterior power and  $S^k = \langle dx_{i_1} \circ \dots \circ dx_{i_k} \rangle$  be the  $k$ -th symmetric power, here

**Table 1.**  $H^2(\mathcal{L}_\delta^+, M_0), L = W_n$

highest weight of $M_0$ and $\hat{M}_0$	$\dim M_0, l = 1, r = 2$ , if otherwise is not mentioned	cocycles $\psi: L_k^+ \wedge L_q^+ \rightarrow \hat{M}_0$ , $\psi(u\partial_i, v\partial_j), 0 \leq k \leq q, k + q = r$
$\pi_1, \lambda = 1,$ $M_0 \cong \Lambda^1$	$n, r = 1$	$\psi_1^W: \partial_i(u) d\partial_j(v) - \partial_j(v) d\partial_i(u)$
$2\pi_1 + \pi_{n-1},$ $\lambda = 2, n > 1,$ $\hat{M}_0 \cong S^2 \otimes L^{\pm 1}$	$(n+2) \binom{n}{2}$	$\psi_2^W: \partial_i(u) v\partial_j - \partial_j(v) u\partial_i$
$\pi_2, \lambda = 1,$ $n > 1, M_0 \cong \Lambda^2$	$\binom{n}{2}$ $l = 2$	$\psi_3^W: d\partial_i(u) \wedge d\partial_j(v);$ $\psi_4^W: d\partial_j(u) \wedge d\partial_i(v)$
$\pi_1 + \pi_2 + \pi_{n-1},$ $\lambda = 2, n > 2$ $\hat{M}_0 \cong \Lambda^1 \otimes \Lambda^2 \otimes L^{\pm 1}$	$(n+2) n^2(n-2)/3$ $l = 2$	$\psi_5^W = \hat{\psi}_2^S:$ $du \wedge d\partial_i(v) \partial_j - dv \wedge d\partial_j(u) \partial_i$ $\psi_6^W:$ $d\partial_i(u) \wedge dv\partial_j - d\partial_j(v) \wedge du\partial_i$
$3\pi_1 + \pi_{n-1},$ $\lambda = 2, n > 1,$ $\hat{M}_0 \cong S^2 \otimes \Lambda^1 \otimes L^{\pm 1}$	$(n+3) \binom{n+1}{3}$	$\psi_7^W: d(\partial_i(u) v) \partial_j - d(\partial_j(v) u) \partial_i$
$2\pi_1 + \pi_2 + 2\pi_{n-1},$ $\lambda = 3, n > 2,$ $\hat{M}_0 \cong S^2 \otimes \Lambda^2 \otimes S^2(L^{\pm 1})$	$\frac{3}{2}(n+4)(n+1) \binom{n+1}{4}$	$\psi_8^W: du \wedge dv\partial_i\partial_j$
$4\pi_1 + \pi_{n-2},$ $\lambda = 2, n > 2,$ $\hat{M}_0 \cong S^3 \otimes \Lambda^1 \otimes \Lambda^2(L^{\pm 1})$	$\binom{n+1}{4} \binom{n+1}{4}$	$\psi_9^W = \hat{\psi}_4^S: d(uv) \partial_i \wedge \partial_j$
$2\pi_2 + \pi_{n-2},$ $\lambda = 2, n > 3,$ $\hat{M}_0 \cong S^2(\Lambda^2) \otimes \Lambda^2(L^{\pm 1})$	$\frac{n+1}{2} \binom{n-2}{2} \binom{n+2}{3}$	$\psi_{10}^W = \hat{\psi}_3^S:$ $\sum_{s,t} (d\partial_s(u) \wedge d\partial_t(v))$ $\circ (dx_s \wedge dx_t) \otimes (\partial_i \wedge \partial_j)$
$\pi_1,$ $\lambda = 0, n = 2$	$2$ $l = 2, r = 3$	$\psi_{11}^W = \hat{\psi}_6^H: d\partial_j(u) \wedge d\partial_i(v)$ $\psi_{12}^W: d\{\partial_i(u), \partial_j(v)\}$
$5\pi_1,$ $\lambda = 2, n = 2$	$6, r = 3$	$\psi_{13}^W: 3\{\overline{pr_1}, \overline{pr_2}\} - 7\overline{pr_1} \wedge \widetilde{\overline{pr_2}}$
$7\pi_1,$ $\lambda = 3, n = 2$	$8, r = 3$	$\psi_{14}^W = \hat{\psi}_7^H: \overline{pr_1} \wedge \overline{pr_2}$
$0, \lambda = -1, n = 1$	$1, r = 2$	$e_0 \wedge e_2 \mapsto 1$
$0, \lambda = -4, n = 1$	$1, r = 5$	$e_2 \wedge e_3 \mapsto 1$
$0, \lambda = -6, n = 1$	$1, r = 7$	$e_2 \wedge e_5 \mapsto 1, e_3 \wedge e_4 \mapsto -3$

**Table 2.**  $H^2(\mathcal{L}_\delta^+, M_0), L = S_{n-1}, n > 2$

highest weight of $M_0$ and $\hat{M}_0$	$\dim M_0, l = 1, r = 2$	cocycles $\psi: L_1^+ \wedge L_1^+ \rightarrow \hat{M}_0$ , its values $\psi(u\partial_i, v\partial_j)$
$\pi_2, M_0 \cong \Lambda^2$	$\binom{n}{2}$	$\psi_1^S: d\partial_j(u) \wedge d\partial_i(v)$
$\pi_1 + \pi_2 + \pi_{n-1},$ $\hat{M}_0 \cong \Lambda^1 \otimes \Lambda^2 \otimes L^{\pm 1}$	$\frac{(n+2)n^2(n-2)}{3}$	$\psi_2^S:$ $du \wedge d\partial_i(v) \partial_j - dv \wedge d\partial_j(u) \partial_i$
$2\pi_1 + \pi_2 + 2\pi_{n-1},$ $\hat{M}_0 \cong S^2 \otimes \Lambda^2 \otimes S^2(L^{\pm 1})$	$\frac{3}{2}(n+4)(n+1) \binom{n+1}{4}$	$\psi_3^S:$ $du \wedge dv\partial_i\partial_j$
$4\pi_1 + \pi_{n-2},$ $\hat{M}_0 \cong S^4 \otimes \Lambda^2(L^{\pm 1})$	$\binom{n+4}{2} \binom{n+1}{4}$	$\psi_4^S:$ $d(uv) \partial_i \wedge \partial_j$
$2\pi_2 + \pi_{n-2}, n > 3,$ $\hat{M}_0 \cong S^2(\Lambda^2) \otimes S^2(L^{\pm 1})$	$\frac{n+1}{2} \binom{n-2}{2} \binom{n+2}{3}$	$\psi_5^S: \sum_{s,t} (d\partial_s(u) \wedge d\partial_t(v))$ $\circ (dx_s \wedge dx_t) \otimes (\partial_i \wedge \partial_j)$

**Table 3.**  $H^1(\mathcal{L}_0^+, M_0)$

$L$	weight	$\dim M_0$	$r$	$l$	cocycles
$W_n$ $n > 1$	$0, \lambda = 1$	1	0	1	$\text{Div}: L_0^+ \rightarrow R(0) \cong \mathbb{C}$
	$\pi_1, \lambda = 1$	$n$	1	1	$\text{Div}: L_1^+ \rightarrow R(\pi_1) \cong (\mathbb{C}_n)_1^+$
	$2\pi_1 + \pi_{n-1}$ $\lambda = 2$	$\binom{n+2}{3}$	1	1	$\overline{pr}_1: L_1^+ \rightarrow R(2\pi_1 + \pi_{n-1}) \cong (S_{n-1})_1^+$
$W_1$	$0, \lambda = 1$	1	0	1	$e_0 \mapsto 1$
	$0, \lambda = 0$	1	1	1	$e_1 \mapsto 1$
	$0, \lambda = -1$	1	2	1	$e_2 \mapsto 1$
$S_{n-1}$ $n > 2$	$2\pi_1 + \pi_{n-1}$	$\binom{n+2}{3}$	1	1	$pr_1: L_1^+ \rightarrow R(2\pi_1 + \pi_{n-1}) \cong L_1^+$
$H_n$	$3\pi_1$	$\binom{2n+2}{3}$	1	1	$pr_1: L_1^+ \rightarrow R(3\pi_1) \cong (\mathbb{C}_{2n})_3^+$
$K_{n+1}$	$3\pi_1, \lambda = 2$	$\binom{n+2}{3}$	1	1	$pr_1: L_1^+ \rightarrow R(3\pi_1) \cong (\mathbb{C}_{2n})_3^+$
	$\pi_1, \lambda = 0$	$2n+1$	1	1	$\partial_0: L_1^+ \rightarrow R(\pi_1) \cong (\mathbb{C}_{2n})_1^+$
	$0, \lambda = 0$	1	0	1	$\partial_0: L_0^+ \rightarrow \mathbb{C}$

**Table 4.**  $H^2(\mathcal{L}_0^+, M_0), L = H_n$

highest weight of $M_0$ and $\hat{M}_0$	$\dim M_0, l = 1, r = 2,$ if otherwise is not mentioned	cocycles $\psi: L_k^+ \wedge L_q^+ \rightarrow \hat{M}_0$ and $\psi(u \wedge v), 0 < k \leq q, k + q = r,$ nonwritten components are zero
$0, M_0 \cong \mathbb{C},$ $\pi_2,$ $\hat{M}_0 \cong \Lambda^2,$ $n > 1$	1 $(n-1)(2n+1)$	$\psi_1^H = \mu$ $\psi_2^H:$ $\sum_{i,j} -d\partial_{-i}\partial_{-j}(u) \wedge d\partial_i\partial_j(v)$ $+ 2d\partial_{-j}\partial_i(u) \wedge d\partial_{-i}\partial_j(v)$ $- d\partial_i\partial_j(u) \wedge d\partial_{-i}\partial_{-j}(v)$
$2\pi_2, n > 1,$ $\hat{M}_0 \cong S^2(\Lambda^2)$	$\frac{2(2n+3)(2n-1)}{3} \binom{n}{2}$	$\psi_3^H:$ $\sum_{i,j} \sum_{s=1}^n (d\partial_i\partial_s(u) \wedge d\partial_j\partial_{-s}(v) - d\partial_i\partial_{-s}(u) \wedge d\partial_j\partial_s(v)) \circ dx_i \wedge dx_j,$
$3\pi_2, n > 1,$ $\hat{M}_0 \cong S^3(\Lambda^2)$	$\frac{(2n+5)(2n+1)(2n-1)n}{9} \binom{n}{2}$	$\psi_4^H:$ $\sum_{i,j,s,t} (d\partial_s\partial_i(u) \wedge d\partial_t\partial_j(v)) \circ (dx_s \wedge dx_t) \circ (dx_i \wedge dx_j),$
$4\pi_1 + \pi_2,$ $\hat{M}_0 \cong S^4 \otimes \Lambda^2,$ $n > 1$	$\frac{(2n+5)(2n+3)(2n+1)}{3} \binom{n+1}{3}$	$\psi_5^H: du \wedge dv$
$\pi_1,$ $M_0 \cong \Lambda^1$	$2n, r = 3$	$\psi_6^H: pr_1\mu(u \wedge v),$
$7\pi_1, n = 1,$ $M_0 \cong (\mathbb{C}_2)_7^+$	$8, r = 3$	$\psi_7^H: uv, u \in L_1^+, v \in L_2^+$
$2\pi_1, n = 1,$ $M_0 \cong (\mathbb{C}_2)_2^+$	$3, r = 4$	$\psi_8^H: pr_2(6\mu + \mu_1)(u \wedge v)$

$i_1, \dots, i_k \in \{1, \dots, n\},$  if  $L = W_n, S_{n-1}$  and  $i_1, \dots, i_k \in \{\pm 1, \dots, \pm n\},$  if  $L = H_n, K_{n+1}.$

It is easy to see that

$$H_2^+(\mathcal{L}_0^+, L_0^+, M_0) \cong C^2(L_1^+, M_0)^{L_0^+}$$

Since the commutators of 0-components of  $W_n^+, S_{n-1}^+$  and

$H_n^+, K_{n+1}^+$  coincide, and the natural imbeddings of 1-components

$$L_1^+(S) = (S_{n-1})_1 \hookrightarrow L_1^+(W) = (W_n^+)_1,$$

$$L_1^+(H) = (H_n)_1^+ \hookrightarrow L_1^+(K) = (K_{n+1})_1$$

are  $[L_0^+, L_0^+]$ -module monomorphisms, we can give

**Table 5.**  $H^2(\mathcal{L}_0^+, L_0^+, M_0), L = K_{n+1}$

highest weight of $M_0$ and $\hat{M}_0$	$\dim M_0, l = 1, r = 2$ , if otherwise is not mentioned	cocycles $\psi$ and its values $\psi(u \wedge v), u \in L_k^+, v \in L_q^+, 0 < k \leq q, k + q = r$
$0, \lambda = -2,$ $M_0 \cong \mathbb{C}$	1	$\psi_1^K = \hat{\psi}_1^H$
$\pi_2, \lambda = -2,$ $\hat{M}_0 \cong \Lambda^2, n > 1$	$(n-1)(2n+1)$ $l = 2$	$\psi_2^K = \hat{\psi}_2^H;$ $\psi_3^K: d\hat{\partial}_0(u) \wedge d\hat{\partial}_0(v)$
$2\pi_2, \lambda = -4,$ $\hat{M}_0 \cong S^2(\Lambda^2), n > 1$	$\frac{2}{3}(2n+3)(2n-1) \binom{n}{2}$	$\psi_4^K = \hat{\psi}_3^H;$
$3\pi_2, \lambda = -6,$ $\hat{M}_0 \cong S^3(\Lambda^2), n > 1$	$\frac{(2n+5)(2n+1)(2n-1)n}{9} \binom{n}{2}$	$\psi_5^K = \hat{\psi}_4^H$
$4\pi_1 + \pi_2, \lambda = 0,$ $\hat{M}_0 \cong S^4(\Lambda^1) \otimes \Lambda^2, n > 1$	$\frac{(2n+5)(2n+3)(2n+1)}{3} \binom{n+1}{3}$	$\psi_6^K = \hat{\psi}_5^H$
$2\pi_1 + \pi_2,$ $\hat{M}_0 \cong S^2 \otimes \Lambda^2,$ $\lambda = -2, n > 1$	$(2n+3)(2n+1) \binom{n}{2}$	$\psi_7^K: \hat{\partial}_0(du \wedge dv)$
$4\pi_1, \lambda = 2,$ $\hat{M}_0 \cong S^4$	$\frac{(2n+3)(2n+1)}{3} \binom{n+1}{2}$	$\psi_8^K: d(u\hat{\partial}_0(v) - v\hat{\partial}_0(u))$
$\pi_1, \lambda = -2,$ $M_0 \cong (\mathbb{C}_2)_1^+,$ $n = 1$	2 $l = 2,$ $r = 3$	$\psi_5^K = \hat{\psi}_6^H$ (prolongation see below) $\psi_9^K_0: x_0^2 \wedge x_0 x_i \mapsto x_i,$
$3\pi_1, \lambda = 0,$ $\hat{M}_0 \cong (\mathbb{C}_2)_3^+$ $n = 1$	4, $r = 3$	$\psi_{11}^K: \hat{\partial}_0(u) \hat{\partial}_0(v) + 3u\hat{\partial}_0^2(v), u \in L_1^+, v \in L_2^+$
$5\pi_1, \lambda = 2,$ $M_0 \cong (\mathbb{C}_2)_5^+$ $n = 1$	6, $r = 3$	$\psi_{12}^K: 3u\hat{\partial}_0(v) + \hat{\partial}_0(u)v$ $u \in L_1^+, v \in L_2^+$
$7\pi_1, \lambda = 0,$ $M_0 \cong (\mathbb{C}_2)_7^+, n = 1$	8, $r = 3$	$\psi_{13}^K = \hat{\psi}_7^H$
Prolongation of $\psi_6^K: \psi_5^K(u \wedge x_0 a) = 5(\sum_{i,j=\pm 1} \text{sgn}(ij) \hat{\partial}_i \hat{\partial}_j(a) \hat{\partial}_{-i} \hat{\partial}_{-j}(u)),$ where $u \in L_1^+, v = x_0 a \in L_2^+.$		

imbeddings

$$H_2^2(\mathcal{L}(S)_0^+, M_0) \hookrightarrow H_2^2(\mathcal{L}(W)_0^+, M_0),$$

$$H_2^2(\mathcal{L}(H)_0^+, M_0) \hookrightarrow H_2^2(\mathcal{L}_0^+, M_0).$$

So, if  $\psi$  is a cocycle from  $Z_2^2(\mathcal{L}(S)_0^+, M_0)$  or  $Z_2^2(\mathcal{L}(H)_0^+, M_0)$ , then it has a trivial prolongation to  $Z_2^2(\mathcal{L}(W)_0^+, M_0)$  or  $Z_2^2(\mathcal{L}(K)_0^+, M_0)$ , correspondingly. Denote it by  $\hat{\psi}$ . This denotation we will reserve also for cocycles from  $Z_r^2(\mathcal{L}(S)_0^+, M_0), Z_r^2(\mathcal{L}(H)_0^+, M_0), r > 2$ , that have prolongations to cocycles from  $Z_r^2(\mathcal{L}(W)_0^+, M_0), Z_r^2(\mathcal{L}(K)_0^+, M_0)$ . Notice that  $S_1 \cong H_1$  and  $sp_1 \cong sl_2$ , but we consider  $S_{n-1}$  from  $n > 2$ , thus in constructing cocycles for  $W_2$  we use cocycles for  $H_1$ .

**6 Main results**

**Theorem 6.1.** Let  $L = W_n, S_{n-1}, H_n, K_{n+1}, Q = L, L^+, (Q, V) = (L, U), (L^+, U), (L^+, U^+)$  and  $M_0$  be irreducible

$\mathcal{L}_0^+$ -module, such that  $\mathcal{L}_1^+ M_0 = 0$ . Then for  $k = 0, 1, 2$ ,

$$H^k(Q, V \otimes M_0) \cong \bigoplus_{s=0}^k H^s(J(L), V) \otimes H^{k-s}(\mathcal{L}_0^+, M_0),$$

$$H^k(\mathcal{L}_0^+, M_0) \cong L_0^+ / [L_0^+, L_0^+] \otimes (H^{k-1}(\mathcal{L}_1^+, \mathbb{C}) \otimes M_0)^{L_0^+} \oplus (H^k(\mathcal{L}_1^+, \mathbb{C}) \otimes M_0)^{L_0^+},$$

$$H^k(J(L), U) \cong \Lambda^k(J(L)), H^k(J(L), U^+) \cong \mathbb{C}^{\delta(L=H_n)}.$$

$H^1(Q, U \otimes M_0)$  has a basis consisting of classes of cocycles  $\delta(M_0 \cong \mathbb{C})Sq_i, i \in I, \delta(L = H_n)\delta(M_0 \cong \mathbb{C})\Delta$ , and  $E\psi(1)$ . In  $H^2(Q, U \otimes M_0)$  one can choose a basis consisting of classes of cocycles  $E(\psi(2)), Sq_i \wedge E(\psi(1)), \delta(L = H_n)\Delta \wedge E(\psi(1)), \delta(M_0 = \mathbb{C})Sq_i \wedge Sq_j$ . Here  $\psi(s)$  denotes the basic cocycles of  $H^s(\mathcal{L}_0^+, M_0), s = 0, 1, 2$ . The space  $H^2(L^+, U^+ \otimes M_0)$  has a basis consisting of classes of cocycles  $E(\psi(2)), \delta(L = H_n)\Delta \wedge E(\psi(1))$ . Nonzero  $H^k(\mathcal{L}_0^+, M_0), k = 1, 2$  see Tables 1–5.

In the tables we use some special denotions:

$L = W_2, \{u, v\}$  denote the Poisson bracket in subalgebra  $H_1$  of  $L, \bar{L}_i$  is the  $L_0$ -submodule of  $L_i$  generated by



derivations without divergence and  $\widetilde{L}_i$  is its additional submodule,  $\overline{pr}_i: L_i \rightarrow \overline{L}_i$ ,  $\widehat{pr}_i: L_i \rightarrow \widehat{L}_i$  are natural projections;

$$L = W_1.e_i = x^{i+1}\partial_i;$$

$L = H_n.pr_i: U \rightarrow U_i$  is natural projection,  $\mu = (\sum_{i=1}^n \partial_{-i} \wedge \partial_i)^3$  is the Moyal cocycle,  $\mu_1$  is the restriction of  $\mu$  to the subalgebra  $\mathcal{L}_2^+$ , prolonged trivially to  $\mathcal{L}_1^+$ , i.e.  $\mu_1(L_1^+, \mathcal{L}_1^+) = 0$ ;

**Corollary 6.2.** *Let  $L = W_n, S_{n-1}, n > 2, H_n, K_{n+1}$ . Then  $H^2(L, L) = H^2(L^+, L^+) = 0$ , except the following cases:  $L = S_n, H_n$ . Moreover,*

$$L = S_2, H^2(L^+, L^+)$$

$$\cong \mathbb{C} \cong \langle u\partial_i \wedge v\partial_j \mapsto \sum_{\{abc\}=\{0,1,2\}} \text{sgn}(\delta_{12}^{abc}) \partial_a \partial_j(u) \partial_b \partial_i(v) \partial_c \rangle,$$

$$L = H_n, Q = L, L^+, H^2(Q, Q) \cong H^1(Q, Q) \wedge H^1(Q, Q) \oplus \langle \mu \rangle, H^1(L, L) \cong H^1(L, U) \cong \langle Sq_i, \Delta, D_\theta \rangle, H^1(L^+, L^+) \cong H^1(L^+, U^+) \cong \langle \Delta, D_\theta \rangle.$$

*Acknowledgements.* The Author is deeply grateful to the Alexander von Humboldt and INTAS foundations for support and to prof. B. Pareigis for hospitality and to prof. A. I. Kostrikin, P. Schauenburg, M. Gerstenhaber, J. Wess, Yu. I. Manin for interest to results.

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